Irrational anyons under an elastic membrane

Claudio Chamon

Collaborators:

Siavosh Behbahani Ami Katz



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$$\mathcal{L} = -\frac{p}{4\pi} \epsilon^{\mu\nu\lambda} a_{\mu} \partial_{\nu} a_{\lambda} + \frac{e}{2\pi} \epsilon^{\mu\nu\lambda} A_{\mu} \partial_{\nu} a_{\lambda}$$

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$$J^{\mu} = \frac{\delta S}{\delta A_{\mu}} = \frac{e}{2\pi} \epsilon^{\mu\nu\lambda} \partial_{\nu} a_{\lambda}$$
$$0 = \frac{\delta S}{\delta a_{\mu}} = -\frac{p}{2\pi} \epsilon^{\mu\nu\lambda} \partial_{\nu} a_{\lambda} + \frac{e}{2\pi} \epsilon^{\mu\nu\lambda} \partial_{\nu} A_{\lambda}$$

$$J^{\mu} = \frac{1}{p} \, \frac{e^2}{h} \, \epsilon^{\mu\nu\lambda} \, \partial_{\nu}A_{\lambda}$$

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$$J_i = \frac{1}{p} \, \frac{e^2}{h} \, \epsilon_{ij} \, E_j$$

$$\rho = \frac{1}{p} \, \frac{e^2}{h} \, \hat{z} \cdot (\vec{\nabla} \times \vec{A})$$

 $Q = \frac{1}{p} e \frac{\Phi}{\phi_0}$



$$J^{\mu} = \frac{1}{p} \, \frac{e^2}{h} \, \epsilon^{\mu\nu\lambda} \, \partial_{\nu}A_{\lambda}$$



Does statistical angle need to be rational?

counter-examples:

. . .

Wilczek, PRL (1982)

Jackiw and Nair, Phys. Lett. B (2000)

Fitzpatrick, Kachru, Kaplan, Katz, and Wacker, arXiv:1205.6816

counter-examples with realization on a lattice:

Chamon, Hou, Jackiw, Mudry, Pi, and Schnyder, PRL (2008)

Ryu, Mudry, Hou, and Chamon, PRB (2009)

Perspective Departure from Chern-Simons Theory

$$a_i = \alpha \epsilon_{ij} \nabla_j \phi$$
 $H = \frac{1}{2m} \psi^{\dagger} (p_i - \alpha \epsilon_{ij} \nabla_j \phi)^2 \psi$

 $b = \epsilon_{ij} \nabla_i a_j = -\alpha \, \nabla^2 \phi$

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$$\oint_{\Phi = -\alpha g/\tau} -e$$

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Behbahani, Chamon, and Katz, arXiv:1307.



$$S = \int dt \, d^2x \left[\frac{\rho}{2} \, \dot{\phi}^2 - \frac{\tau}{2} \, (\nabla \phi)^2 - \frac{\kappa}{2} \, (\nabla^2 \phi)^2 \right. \\ \left. + \psi^{\dagger} i \partial_t \psi - \frac{1}{2m} \left| (\nabla_i - i\alpha \, \epsilon_{ij} \nabla_j \phi) \, \psi \right|^2 - g \, \phi \, \psi^{\dagger} \psi \right]$$



$$H_{\psi} = \frac{1}{2m} \left(\vec{p} + e\vec{A} \right)^2 - eA_0$$



$$A_{i}(X) = \int d^{2}X' \frac{\sigma_{m} \epsilon_{ij} (X - X')_{j}}{\left[(X - X')^{2} + (\delta + \phi (X'))^{2} \right]^{3/2}},$$

 $\vec{A} = \vec{A}^{BG} + \vec{A^{\phi}}$

 $A_i^{\phi}(X) \approx 2\pi \,\sigma_m \,\epsilon_{ij} \nabla_j \phi(X) \operatorname{sgn}(\delta) \quad \longrightarrow \quad \alpha = -2\pi \,\sigma_m \, e \, \operatorname{sgn}(\delta)$ $A_0^{\phi}(X) \approx -2\pi \,\sigma_e \, \phi(X) \operatorname{sgn}(\delta) \quad \longrightarrow \quad g = 2\pi \,\sigma_e \, e \, \operatorname{sgn}(\delta)$

$$\theta = -2\pi^2 e^2 \ \frac{\sigma_e \sigma_m}{\tau}$$



Chamon, Freed, Kivelson, Sondhi, Wen, PRB (1997)





Chamon, Freed, Kivelson, Sondhi, Wen, PRB (1997)

$$\Delta B^* = \Phi^* / A \qquad \Phi^* = \frac{e}{e^*} \Phi_0$$
$$\Delta B = \Phi_0 / A \qquad \Phi^* = \frac{e}{e^*} \Phi_0$$















Proposed experimental set-up to probe anyons under a membrane





Proposed experimental set-up to probe anyons under a membrane

membrane layer





$$\nabla^2 \phi(\vec{r}) = \frac{g}{\tau} \,\rho(\vec{r}) - \frac{\alpha}{\tau} \,\hat{z} \cdot \vec{\nabla} \times \vec{J}(\vec{r})$$

$$B = \frac{\alpha}{e} \, \nabla^2 \phi$$

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$$B = \frac{\alpha}{e} \nabla^2 \phi \qquad Q = -en$$

$$P = -\frac{\alpha g}{\tau} \frac{Q}{e^2}$$

$$= -\frac{\alpha g}{\tau} \frac{Q}{e^2}$$

Proposed experimental set-up to probe anyons under a membrane

$$G(\Phi_{\text{total}}/\Phi_{0})$$

$$\Phi_{\text{total}} = \Phi + \Phi_{\text{ext}}$$

$$2\pi \Delta \Phi/\Phi_{0} = -(\alpha g/\tau) \Delta Q/e$$

$$\theta = \frac{\alpha g}{2\tau} = \pi \frac{\Delta \Phi_{\text{ext}} / \Phi_0}{\Delta Q / e} = \frac{1}{2} \frac{e^2}{C} \frac{\Delta \Phi_{\text{ext}}}{\Delta V}$$

Possible experimental realization



electron gas layer: metal



Magnetic impurities spaced about 1 nm away, with each impurity having a magnetic moment of order μ_B , yields a magnetic density of order $\sigma_m \sim 3 \times 10^{-3} \, eV$. $R\sim 3\mu m$

 $\Delta R = R_D - R \sim 0.1 \, R \sim 300 \, nm$

 $\phi_{\rm max} \sim 300 \ nm$

Possible experimental realization

$$-\frac{Q}{e}\frac{\theta}{\pi} = 2\pi\sigma_m e \phi_{\max} \frac{1}{\ln R_D/R}$$
$$\approx 2\pi\sigma_m e \phi_{\max} \frac{R}{\Delta R}$$
$$\sim 3 \times 10^{-2}$$



 $N_e \sim 1$

log potential

$$S = \int dt \, d^2x \left[\frac{\rho}{2} \, \dot{\phi}^2 - \frac{\tau}{2} \, (\nabla \phi)^2 - \frac{\kappa}{2} \, (\nabla^2 \phi)^2 \right. \\ \left. + \psi^{\dagger} i \partial_t \psi - \frac{1}{2m} \left| (\nabla_i - i\alpha \, \epsilon_{ij} \nabla_j \phi) \, \psi \right|^2 - g \, \phi \, \psi^{\dagger} \psi \right]$$

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no log potential

$$\nabla^2 \phi \ \psi^{\dagger} \sigma_z \psi \qquad \nabla^2 \phi \ \psi^{\dagger} \psi$$

Zeeman

quadrupole

log potential

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no log potential

$$\nabla^2 \phi \ \psi^{\dagger} \sigma_z \psi \qquad \qquad \nabla^2 \phi \ \psi^{\dagger} \psi$$

Zeeman

quadrupole



metal: screening of interactions!

Amusement: log-interaction allows the mesoscopic study of "anyon stars"



vary θ

Conclusions

Anyon statistics mediated by electromagnetism and phonons

Statistical angle is a *continuous function* of membrane properties

- magnetic dipolar density
- electric charge density
- tension

Proposed physical realization with graphene as membrane

Proposed experiment that does not suffer from "area" problem