

# Spectral properties of Laplacians on non-compact manifolds with general ends

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Inverse Problems and Related Topics



# Outline

- 1 Introduction
  - Manifolds with ends

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# Works I

-  H. Isozaki, Y. Kurylev and M. Lassas,  
Forward and inverse scattering problem on manifolds with asymptotically hyperbolic ends,  
*J. Funct. Anal.* 258 (2010), 2010-2118.
-  H. Isozaki, Y. Kurylev and M. Lassas,  
Conic singularities, generalized scattering matrix, and inverse scattering on asymptotically hyperbolic surfaces,  
to appear in *J. Reine angew. Math.*

# Works II



H. Isozaki, Y. Kurylev and M. Lassas,  
Spectral theory and inverse problems on asymptotically  
hyperbolic orbifolds,  
*arXiv : 1312.0421*

We consider a connected, non-compact manifold (or orbifold) of the form

$$\mathcal{M} = \mathcal{K} \cup \mathcal{M}_1 \cup \cdots \cup \mathcal{M}_{N+N'}$$

where,  $\mathcal{K}$  = relatively compact,

$$\mathcal{M}_i \simeq (0, \infty) \times M_i, \quad (\text{diffeomorphic}),$$

$M_i$  = compact manifold (orbifold) of dim  $n - 1$  equipped with the metric  $h_M(x, dx)$ .

$\mathcal{M}_i$  is equipped with the metric

$$ds^2 = (dr)^2 + \rho_i(r)^2 h_i(r, x, dx)$$

$$h_i(r, x, dx) - h_{M_i}(x, dx) = O(r^{-\gamma}), \quad \gamma \geq 0.$$

# Typical examples

- (1)  $\rho(r) = e^{c_0 r}$ ,  $c_0 > 0$ , hyperbolic regular end
- (2)  $\rho(r) = r$  Euclidean (or conical)
- (3)  $\rho(r) = 1$  cylindrical end (waveguide)
- (4)  $\rho(r) = e^{c_0 r}$ ,  $c_0 < 0$ , hyperbolic cusp

## We know that

- For (1), (2) No embedded eigenvalues in the continuous spectrum.
- For (3), (4)  $\exists$  embedded eigenvalues
- For (1), (2), (3) One can reconstruct  $\mathcal{M}$  from the **physical S-matrix** for all energies associated with one end
- For (4), One can reconstruct  $\mathcal{M}$  from the **generalized S-matrix** for all energies associated with one end



We deal with the metric of the form

$$\rho(r) \sim \begin{cases} \exp\left(c_0 r + \frac{\beta}{\alpha} r^\alpha\right), & 0 \leq \alpha < 1, \\ r^\beta, & \end{cases}$$

where

- For the regular end,  $c_0 > 0$  or  $c_0 = 0, \beta > 0$ ,
- For the cusp,  $c_0 < 0$ , or  $c_0 = 0, \beta < 0$ .

Sometimes, it is more convenient to state the assumption in the form

$$\rho(r)^{-1} \in \mathcal{S}^{-\beta}.$$

In this setting, the exponentially growing metric corresponds to the case  $\beta = \infty$ .

# Outline

- Assumptions
  - Instability of the short-range assumption

2 Rellich-Vekua theorem

3 Laplacian on  $\mathcal{M}$

4 Spectral representation and S-matrix

5 Inverse scattering

# Perturbed warped product metric

Consider the following metric on  $(0, \infty) \times M$ ,

$$ds^2 = (dr)^2 + g_M(r, x, dx),$$

where  $g_M(r, x, dx) = g_{M,ij}(r, x)dx^i dx^j$  is a metric on  $M$  depending smoothly on  $r > 0$ . Let

$$g = g(r, x) = \det(g_{M,ij}(r, x)).$$

Define

$$f(r, x) \in S^\kappa \iff \partial_r^m \partial_x^\alpha f(r, x) = O(r^{\kappa-m}), \quad \forall m, \alpha.$$

## The 1st assumption

$$(A-1) \quad \frac{g'}{4g} - \frac{(n-1)c_0}{2} - \frac{(n-1)\beta}{2} r^{\alpha-1} \in S^{-1-\epsilon}, \quad \epsilon > 0.$$

$$0 \leq \alpha < 1, \quad \epsilon > 0,$$

$$\beta \neq 0, \text{ if } c_0 = 0.$$

Integrating

$$\frac{g'}{4g} - \frac{(n-1)c_0}{2} - \frac{(n-1)\beta}{2} r^{\alpha-1} = O(r^{-1-\epsilon}),$$

we have

$$g = \rho(r)^{2(n-1)} O(1),$$

where

$$\rho(r) = \begin{cases} \exp\left(c_0 r + \frac{\beta}{\alpha} r^\alpha\right), & 0 < \alpha < 1, \\ \exp(c_0 r) r^\beta, & \alpha = 0. \end{cases}$$

We can then rewrite  $g_M(r, x, dx)$  as

$$g_M(r, x, dx) = \rho(r)^2 h(r, x, dx),$$

where  $h(r, x, dx)$  is bounded in  $r$ .

So, our metric has the form

$$ds^2 = (dr)^2 + \rho(r)^2 h(r, x, dx).$$

## The 2nd assumption

(A-2) There exists a smooth metric  $h_M(x, dx)$  on  $M$  such that

$$h(r, x, dx) - h_M(x, dx) \in \mathcal{S}^{-\gamma}, \quad \gamma > 0.$$

We say that

- $\mathcal{M}$  has a **regular infinity** if

$$\text{either } c_0 > 0, \quad \text{or } c_0 = 0, \quad \beta > 0.$$

- $\mathcal{M}$  has a **cusp** if

$$\text{either } c_0 < 0, \quad \text{or } c_0 = 0, \quad \beta < 0.$$

## The 3rd assumption

(A-3)  $\mathcal{M}_1, \dots, \mathcal{M}_N$  have regular infinities, and  $\mathcal{M}_{N+1}, \dots, \mathcal{M}_{N+N'}$  have cusp.

## The perturbation term

$$h(r, x, dx) - h_M(x, dx) \in \mathcal{S}^{-\gamma}$$

is said to be **short-range** if  $\gamma > 1$ , and **long-range** if  $\gamma \leq 1$ .

Usually, the latter is more complicated than the former.

It seems that there exist **threshold growth orders** of the metric :

Assuming that  $\rho(r) = r^\beta$ , they are (think of

$$x_{n+1} = (x_1^2 + \cdots + x_n^2)^{\beta/2})$$

- $\beta = 1$  : conic surface
- $\beta = 1/2$  : parabola
- $\beta = 1/3$  : ?

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# Instability of the short-range assumption

Consider the following metric with **a cross term**

$$ds^2 = a(t, z)(dt)^2 + 2w(t)b_{ij}(t, z)dtdz^i + w(t)^2c_{ij}(t, z)dz^i dz^j.$$

Assume

$$w(t)^{-1} \in \mathcal{S}^{-\kappa}, \quad a(t, z) - 1 \in \mathcal{S}^{-\lambda},$$

$$b_{ij}(t, z) \in \mathcal{S}^{-\mu}, \quad c_{ij}(t, z) - h_{ij}(t, z) \in \mathcal{S}^{-\nu},$$

with the condition

$$\kappa > 1/2, \quad \lambda > 1, \quad \kappa + \mu > 1, \quad \nu > 0.$$

Note that  $\kappa$  corresponds to the **volume growth** of the manifold.

## Theorem

One can transform the metric with a cross term

$$ds^2 = a(t, z)(dt)^2 + 2w(t)b_i(t, z)dtdz^i + w(t)^2c_{ij}(t, z)dz^i dz^j.$$

into the **perturbed warped product** form

$$ds^2 = (dr)^2 + w(r)^2\bar{h}(r, x, dx),$$

where  $\bar{h}(r, x, dx)$  is an  $r$ -dependent metric on  $M$  satisfying

$$\bar{h}(r, x, dx) - h(z(x), dx) \in \mathcal{S}^{-\min\{\nu, \epsilon_0\}},$$

$$\epsilon_0 = \min\{\lambda, \kappa + \mu, 2\kappa\} - 1.$$

Therefore, the metric with cross term can be transformed to the

perturbed warped product if

$$w(t) \sim \exp\left(c_0 t + \frac{\beta}{\alpha} t^\alpha\right), \quad \text{or } t^\beta, \quad \text{with } \beta > 1/2.$$

However, even if  $\nu$  is large, the resulting metric  $(dr)^2 + w(r)^2 \bar{h}(r, x, dx)$  is

- a **short-range** perturbed metric of  $(dr)^2 + \rho(r)^2 h_M(x, dx)$  only when  $\beta > 1$ ,
- a **long-range** perturbed metric if  $1/2 < \beta \leq 1$ .

Note that the case  $\beta = 1$  corresponds to the standard asymptotically Euclidean metric (see [Bouclet, 2012]).

## Rellich-Vekua theorem

Consider  $\mathcal{M} = (0, \infty) \times M$  with metric

$$ds^2 = (dr)^2 + \rho(r)^2 h(r, x, dx),$$

$$h(r, x, dx) - h_M(x, dx) \in S^{-\epsilon}, \quad \epsilon > 0.$$

Assume that

(B-1) *There exist constants  $c_0, \epsilon_0, \delta_0, \epsilon$  such that*

$$\frac{\rho'}{\rho} - c_0 \in S^{-\epsilon_0}, \quad (r\partial_r + \delta_0)\rho^{-1} \leq 0$$

*and satisfying either (i) or (ii) :*

(i)  $c_0 \geq 0, \delta_0 > 1/3, \epsilon_0 > 0, \epsilon > 0.$

(ii)  $c_0 = 0, \delta_0 > 0, \epsilon_0 = 1, \epsilon > 0.$

We fix a point  $p_0$  in  $\mathcal{M}$ , and let

$$S(r) = \{p \in \mathcal{M}; \text{dist}(p, p_0) = r\}.$$

### Theorem

Suppose there exist constants  $R > 0$ ,  $\lambda > ((n-1)c_0/2)^2$  and  $u \in H_{loc}^2(\mathcal{M})$  such that

$$(-\Delta_{\mathcal{M}} - \lambda)u = 0, \quad \text{for } r > R,$$

$$\liminf_{r \rightarrow \infty} r^\gamma \int_{S(r)} \left( \left| \frac{\partial u}{\partial r} \right|^2 + |u|^2 \right) dS(r) = 0, \quad \exists \gamma > 0.$$

Then  $u = 0$  for  $r > R$ .

This theorem covers all of the cases :

$$\rho(r) = \begin{cases} \exp\left(c_0 r + \frac{\beta}{\alpha} r^\alpha\right), & 0 < \alpha < 1, \\ \exp(c_0 r) r^\beta, & \alpha = 0. \end{cases}$$

To prove this theorem, we consider an abstract differential equation

$$-u''(t) + B(t)u(t) + V(t)u(t) - \lambda u(t) = 0, \quad t > 0$$

for an Hilbert space - valued functions, and apply the classical method of T. Kato ([1959], CPAM), or Eidus ([1969], Russ. Math. Survey).

Here,  $B(t)$  (corresponding to  $-\rho(r)^{-2}\Delta_M$ ) is a non-negative self-adjoint operator having the property

$$t \frac{dB(t)}{dt} + \delta B(t) \leq Ct^{-\epsilon}, \quad \delta > 0.$$

## Basic spectral properties

Consider the case of one end :  $\mathcal{M} = (0, \infty) \times M$ , on which the metric is

$$ds^2 = (dr)^2 + \rho(r)^2 h(r, x, dx).$$

Then, the Laplacian is

$$-\Delta_{\mathcal{M}} = -\partial_r^2 - \frac{g'}{2g} \partial_r + \rho^{-2} B(r),$$

$$g = \rho^{2(n-1)} h, \quad h = h(r, x) = \det(h_{ij}),$$

$$B(r) = -\frac{1}{\sqrt{h}} \partial_{x_i} \left( \sqrt{h} h^{ij} \partial_{x_j} \right).$$

We put

$$\mathbf{h}(r) = L^2(M; \sqrt{h(r, x)} dx),$$

$$L^{2,s}(\mathcal{M}) \ni f \iff \|f\|_s^2 = \int_0^\infty (1+r)^{2s} \|f(r)\|_{\mathbf{h}(r)}^2 \rho(r)^{n-1} dr < \infty,$$

$$\mathcal{B} \ni f \iff \sum_{j=1}^\infty 2^{j/2} \|f\|_{L^2(I_j)} < \infty,$$

$$I_0 = (0, 1), \quad I_j = (2^{j-1}, 2^j), \quad (j \geq 1),$$

$$\mathcal{B}^* \ni v \iff \sup_{R>1} \frac{1}{R} \int_0^R \|v(r)\|_{\mathbf{h}(r)}^2 \rho(r)^{n-1} dr < \infty,$$

$$\mathcal{B}_0^* \ni v \iff \lim_{R \rightarrow \infty} \frac{1}{R} \int_0^R \|v\|_{\mathbf{h}(r)}^2 \rho(r)^{n-1} dr = 0.$$



We then have the following inclusion relations

$$L^{2,s} \subset \mathcal{B} \subset L^{2,1/2} \subset L^2 \subset L^{2,-1/2} \subset \mathcal{B}^* \subset L^{2,-s}, \quad s > 1/2.$$

We return to our manifold

$$\mathcal{M} = \mathcal{K} \cup \mathcal{M}_1 \cdots \mathcal{M}_{N+N'}.$$

Take a partition of unity  $\{\chi_j\}_{j=0}^{N+N'}$ , and define the norms on  $\mathcal{M}$ , e.g.

$$\|f\|_s = \|\chi_0 f\|_{L^2(\mathcal{M})} + \sum_{j=1}^{N+N'} \|\chi_j f\|_{L^{2,s}(\mathcal{M}_j)}.$$

Recall that on each end  $\mathcal{M}_i = (0, \infty) \times M_i$ ,

$$ds^2 = (dr)^2 + \rho_i(r)^2 h_i(r, x, dx),$$

$$\rho_i(r) = \begin{cases} \exp(c_{0,i}r + \frac{\beta_i}{\alpha_i}r^{\alpha_i}), & 0 < \alpha_i < 1, \\ \exp(c_{0,i}r)r^{\beta_i}, & \alpha_i = 0, \end{cases}$$

$$h_i(r, x, dx) - h_{M_i}(x, dx) \in \mathcal{S}^{-\gamma_i},$$

where

for **regular ends**  $1 \leq i \leq N$ ,

either  $c_{0,i} > 0$ , or  $c_{0,i} = 0, \beta_i > 0$ ,

for **cusp ends**  $N + 1 \leq i \leq N + N'$ ,

either  $c_{0,i} < 0$ , or  $c_{0,i} = 0, \beta_i < 0$ .

Let

$$H = -\Delta_{\mathcal{M}} \quad \text{on } \mathcal{M},$$

$$R(z) = (H - z)^{-1}.$$

Let

$$E_{0,i} = ((n-1)c_{0,i}/2)^2, \quad \mathcal{T} = \{E_{0,1}, \dots, E_{0,N+N'}\}.$$

$$E_0 = \min_{1 \leq i \leq N+N'} E_{0,i}, \quad E_{0,reg} = \min_{1 \leq i \leq N} E_{0,i}.$$

### Lemma

- (1)  $\sigma_e(H) = [E_0, \infty)$ .
- (2)  $\sigma_p(H) \cap ((E_{0,reg}, \infty) \setminus \mathcal{T}) = \emptyset$ .
- (3) Suppose  $E_0 < E_{0,reg}$ . Then, the eigenvalues in  $(E_0, E_{0,reg}) \setminus \mathcal{T}$  are of finite multiplicities with possible accumulation points at  $\mathcal{T}$ .
- (4) If all ends are cusp, the eigenvalues in  $(E_0, \infty) \setminus \mathcal{T}$  are of finite multiplicities with possible accumulation points at  $\mathcal{T}$  and infinity.

Take any compact interval

$$I \subset (E_0, \infty) \setminus \left( \sigma_p(H) \cup \{E_{0,1}, \dots, E_{0,N+N'}\} \right).$$

Assume long-range, i.e. for  $h_j(r, x, dx) - h_{M_j}(x, dx) \in \mathcal{S}^{-\gamma_j}$ ,

$$\gamma_j > 0, \quad \forall j = 1, \dots, N + N'.$$

### Theorem

For any  $f \in \mathcal{B}$  and  $\lambda \in I$ , there exists a weak  $*$ -limit  
 $\lim_{\epsilon \rightarrow 0} R(\lambda \pm i\epsilon)f$ , i.e. for  $f, g \in \mathcal{B}$

$$\lim_{\epsilon \rightarrow 0} (R(\lambda \pm i\epsilon)f, g) = (R(\lambda \pm i0)f, g).$$

Moreover

$$\|R(\lambda \pm i0)f\|_{\mathcal{B}^*} \leq C \|f\|_{\mathcal{B}}, \quad \lambda \in I.$$

## The representation space

Let

$$\mathbf{h}_j = L^2(M_j).$$

$$\mathbf{h}_{j,\infty} = \begin{cases} \mathbf{h}_j, & 1 \leq j \leq N, \\ \mathbf{C}, & N+1 \leq j \leq N+N', \end{cases}$$

$$\mathbf{h}_\infty = \bigoplus_{j=1}^{N+N'} \mathbf{h}_{j,\infty},$$

$$\widehat{\mathcal{H}} = \bigoplus_{j=1}^{N+N'} L^2((E_{0,j}, \infty); \mathbf{h}_{\infty,j}; d\lambda).$$

Let  $\mathcal{H}_{ac}(H)$  be the absolutely continuous subspace for  $H$ , i.e.

$$\mathcal{H}_{ac}(H) \ni f \iff d(E_H(\lambda)f, f) \text{ is absolutely continuous.}$$

where  $E_H(\lambda)$  is the spectral decomposition of  $H$ .

We assume either

$$\beta_i \geq 1, \gamma_i > 0,$$

or

$$\beta_i > 0, \gamma_i > 1.$$

As in the case of  $\mathbf{R}^n$ , there exists a **spectral representation** (or **generalized Fourier transformation**) associated with  $H$ . It is first defined on  $\mathcal{H}$ ,

$$\mathcal{F}^{(\pm)}(\lambda) = (\mathcal{F}_1^{(\pm)}(\lambda), \dots, \mathcal{F}_{N+N'}^{(\pm)}(\lambda)) \in \mathbf{B}(\mathcal{B}; \mathbf{h}_\infty),$$

(to be explained later),  
and then extended to  $L^2(\mathcal{M})$  by the formula

$$(\mathcal{F}^{(\pm)}f)(\lambda) = \mathcal{F}^{(\pm)}(\lambda)f.$$

## Theorem

There exists a unitary operator  $\mathcal{F}^{(\pm)} : \mathcal{H}_{ac}(H) \rightarrow \widehat{H}$  having the following properties.

- (1)  $(\mathcal{F}^{(\pm)} f)(\lambda) = \mathcal{F}^{(\pm)}(\lambda) f, \forall f \in \mathcal{B}$ .
- (2)  $(\mathcal{F}^{(\pm)} H f)(\lambda) = \lambda (\mathcal{F}^{(\pm)} f)(\lambda), \forall f \in D(H)$ .
- (3)  $\mathcal{F}^{(\pm)}(\lambda)^* \in \mathbf{B}(\mathbf{h}_\infty; \mathcal{B}^*)$ , and

$$(H - \lambda) \mathcal{F}^{(\pm)}(\lambda)^* = 0.$$

(4) For any  $f \in \mathcal{H}_{ac}(H)$ , the inversion formula holds

$$f = \sum_{j=1}^{N+N'} \int_{E_{0,j}}^{\infty} \mathcal{F}_j^{(\pm)}(\lambda)^* (\mathcal{F}_j^{(\pm)} f)(\lambda) d\lambda.$$

## Asymptotic expansion of the resolvent at infinity

Each operator  $\mathcal{F}_j^{(\pm)}(\lambda)$  is constructed as follows.  
Note that  $-\Delta_{h_j}$  has

- eigenvalues  $B_\ell^{(j)}$ ,  $\ell = 0, 1, 2, \dots$ ,
- eigenprojections  $P_\ell^{(j)}$ ,  $\ell = 0, 1, 2, \dots$ .

Then for  $1 \leq j \leq N$ ,

$$\mathcal{F}_j^{(\pm)}(\lambda) = \sum_{\ell=0}^{\infty} \mathcal{F}_{j,\ell}^{(\pm)}(\lambda) \otimes P_\ell^{(j)},$$

for  $N + 1 \leq j \leq N + N'$ ,

$$\mathcal{F}_j^{(\pm)}(\lambda) = \mathcal{F}_{j,0}^{(\pm)}(\lambda).$$

Here  $\mathcal{F}_{j,\ell}^{(\pm)}(\lambda)$  are related to the asymptotic behavior of the  
**resolvent at infinity**:



For  $f \in \mathcal{B}$ ,  $R(\lambda \pm i0)f$  behaves like as  $r \rightarrow \infty$

- on  $\mathcal{M}_j$ ,  $1 \leq j \leq N$ ,

$$\sum_{\ell \geq 0} c_j(\lambda) \rho_j(\lambda)^{-(n-1)/2} e^{i\varphi_{j,\ell}^{(\pm)}(\lambda,r)} \mathcal{F}_{j,\ell}^{(\pm)}(\lambda) \otimes P_\ell^{(j)} f,$$

- on  $\mathcal{M}_j$ ,  $N+1 \leq j \leq N+N'$ ,

$$c_j(\lambda) \rho_j(\lambda)^{-(n-1)/2} e^{i\varphi_{j,0}^{(\pm)}(\lambda,r)} \mathcal{F}_{j,0}^{(\pm)}(\lambda) \otimes P_0^{(j)} f,$$

$$\varphi_{j,\ell}^{(\pm)}(\lambda, r) \sim \pm r \sqrt{\lambda - \frac{(n-1)^2 c_{0,j}^2}{4}}.$$

# Helmholtz equation

Let

$$\mathbf{h}_\infty(\lambda) = \bigoplus_{j=1}^{N+N'} \chi_{(E_{0,j,\infty})}(\lambda) \mathbf{h}_{\infty,j}.$$

## Theorem

$$\{u \in \mathcal{B}^* ; (H - \lambda)u = 0\} = \mathcal{F}^{(\pm)}(\lambda)^* \mathbf{h}_\infty(\lambda).$$

## S-matrix

Take  $\forall \psi^{(in)} = (\psi_1^{(in)}, \dots, \psi_{N+N'}^{(in)}) \in \mathbf{h}_\infty(\lambda)$ , and let

$$\psi_{j,\ell}^{(in)} = P_\ell^{(j)} \psi_j^{(in)}.$$

Then  $\exists ! u \in \mathcal{B}^*$  and  $\exists ! \psi^{(out)} \in \mathbf{h}_\infty(\lambda)$  such that  $(H - \lambda)u = 0$ , which behaves as follows :

(1) on the regular ends  $\mathcal{M}_j$ ,  $1 \leq j \leq N$ ,

$$u \simeq \sum_{\ell=0}^{\infty} \omega_-^{(j,\ell)}(\lambda) \rho_j(\lambda)^{-(n-1)/2} e^{-i\varphi_j(\lambda,r)} \psi_{j,\ell}^{(in)},$$

$$- \sum_{\ell=0}^{\infty} \omega_+^{(j,\ell)}(\lambda) \rho_j(r)^{-(n-1)/2} e^{i\varphi_j(\lambda,r)} \psi_{j,\ell}^{(out)}$$

(2) on the cusp ends  $\mathcal{M}_j$ ,  $N + 1 \leq j \leq N + N'$

$$u \simeq \omega_-^{(j)}(\lambda) \rho_j(\lambda)^{-(n-1)/2} e^{-i\varphi_j(\lambda, r)} \psi_j^{(in)}$$

$$- \omega_+^{(j)}(\lambda) \rho_j(r)^{-(n-1)/2} e^{i\varphi_j(\lambda, r)} \psi_j^{(out)}$$

The **S-matrix** is then defined by

$$\widehat{S}(\lambda) : \mathbf{h}_\infty(\lambda) \ni \psi^{(in)} \rightarrow \psi^{(out)} \in \mathbf{h}_\infty(\lambda)$$

which is unitary.

# Inverse Scattering

## Theorem (Inverse scattering from regular ends)

Given two manifolds  $\mathcal{M}^{(1)}, \mathcal{M}^{(2)}$ , let  $\mathcal{M}_1^{(1)}, \mathcal{M}_1^{(2)}$  be regular ends, and assume that

$$\widehat{\mathcal{S}}_{11}^{(1)}(\lambda) = \widehat{\mathcal{S}}_{11}^{(2)}(\lambda), \quad \forall \lambda,$$

moreover  $\mathcal{M}_1^{(11)}$  and  $\mathcal{M}_1^{(2)}$  are isometric. Then  $\mathcal{M}^{(1)}$  and  $\mathcal{M}^{(2)}$  are isometric.

The physical S-matrix for the cusp end does not have sufficient information to recover the whole manifold.

## Generalized S-matrix

Assume that for the cusp ends  $\mathcal{M}_j$ ,  $j = N + 1, \dots, N + N'$ , the metric has the form

$$ds^2 \Big|_{\mathcal{M}_j} = (dr)^2 + \rho_j(r)^2 h_{M_j}(x, dx),$$

$$\rho_j(r) = \begin{cases} \exp\left(c_{0,j}r + \frac{\beta_j}{\alpha_j}r^{\alpha_j}\right), & 0 < \alpha_j < 1, \\ \exp(c_{0,j}r)r^{\beta_j}, & \alpha_j = 0, \end{cases}$$

where  $c_{0,j} < 0$ , or  $c_{0,j} = 0$ ,  $\beta_j < 0$ .

Let  $0 = B_0^{(j)} \leq B_1^{(j)} \leq B_2^{(j)} \leq \dots$  be the eigenvalues of  $-\Delta_{M_j}$  with normalized eigenvectors  $\phi_\ell^{(j)}$ ,  $\ell = 0, 1, 2, \dots$ . We put

$$\Phi_j(r) = \int_1^r \left( \frac{B_\ell^{(j)}}{\rho_j^2} - \lambda + \frac{n^2 - 2n}{4} \left( \frac{\rho_j'}{\rho_j} \right)^2 \right)^{1/2} dt.$$

Note

$$\Phi_j(r) \rightarrow \infty, \quad \text{as } r \rightarrow \infty.$$

Then, there exist solutions of the equation

$$-u'' - \frac{(n-1)\rho_j'}{\rho_j} u' + \left( \frac{B_\ell^{(j)}}{\rho_j^2} - \lambda \right) u = 0,$$

which behave like

$$u_{\ell,\pm}^{(j)} \sim \rho_j(r)^{-(n-1)/2} e^{\pm \Phi_j(r)}, \quad r \rightarrow \infty.$$

Take any solution  $u$  of the equation

$$(-\Delta_{\mathcal{M}} - \lambda)u = 0, \quad \text{on } \mathcal{M}_j, \quad j = N + 1, \dots, N + N'.$$

Expanding it by  $\phi_{\ell}^{(j)}$ , we have

$$(u, \phi_{\ell}^{(j)})_{L^2(\mathcal{M}_j)} = a_{\ell}^{(j)} u_{\ell,+}^{(j)}(r) + b_{\ell}^{(j)} u_{\ell,-}^{(j)}(r).$$

Here, we introduce two sequences  $\mathbf{A}_{\pm}^{(j)}$  :

$$\mathbf{A}_{\pm}^{(j)} \ni \{c_{\ell,\pm}\}_{\ell=0}^{\infty} \iff \sum_{\ell=0}^{\infty} |c_{\ell,\pm}|^2 |u_{\ell,\pm}^{(j)}(r)|^2 < \infty, \quad \forall r > 1.$$

Then,  $c_{\ell,\pm}^{(j)}$  behaves like, roughly,

$$\mathbf{A}_{\pm}^{(j)} \ni \{c_{\ell,\pm}\}_{\ell=0}^{\infty} \iff \sum_{\ell=0}^{\infty} |c_{\ell,\pm}|^2 e^{\pm 2B_{\ell}^{(j)} \Phi_j(r)} < \infty, \quad \forall r > 1.$$



So,  $\mathbf{A}_+^{(j)}$  is the space of super-exponentially decaying sequences, and  $\mathbf{A}_-^{(j)}$  is the space of super-exponentially growing sequences.

### Lemma

If  $\{a_\ell^{(j)}\} \in \mathbf{A}_+^{(j)}$ , then  $\{b_\ell^{(j)}\} \in \mathbf{A}_-^{(j)}$ .

For the cusp end, we define the generalized incoming data and outgoing data by

$$\Psi_j^{(in)} = \sum_{\ell=0}^{\infty} a_\ell^{(j)} u_{\ell,+}^{(j)}(r) \phi_\ell^{(j)}(x), \quad \{a_\ell^{(j)}\}_{\ell=0}^{\infty} \in \mathbf{A}_+^{(j)},$$

(which is exponentially growing as  $r \rightarrow \infty$ ),

$$\Psi_j^{(out)} = \sum_{\ell=0}^{\infty} b_{\ell}^{(j)} u_{\ell,-}^{(j)}(r) \phi_{\ell}^{(j)}(x), \quad \{b_{\ell}^{(j)}\}_{\ell=0}^{\infty} \in \mathbf{A}_{-}^{(j)},$$

(which is exponentially decaying as  $r \rightarrow \infty$ ).

For the regular end, we put

$$\Psi_j^{(in)} = \sum_{\ell=0}^{\infty} \omega_{-}^{(j)}(\lambda, B_{\ell}^{(j)}) \Psi_j^{(-)}(\lambda, r; B_{\ell}^{(j)}) \psi_{j,\ell}^{(in)},$$

$$\Psi_j^{(out)} = \sum_{\ell=0}^{\infty} \omega_{-}^{(j)}(\lambda, B_{\ell}^{(j)}) \Psi_j^{(-)}(\lambda, r; B_{\ell}^{(j)}) \psi_{j,\ell}^{(out)}.$$

## Theorem

For any incoming data  $\psi_j^{(in)} \in L^2(\mathcal{M}_j)$  from regular ends and generalized incoming data from cusp ends, there exist a solution  $u$  of the equation  $(-\Delta_{\mathcal{M}} - \lambda)u = 0$  and the outgoing data such that

$$u - \sum_{j=N+1}^{N+N'} \Psi_j^{(in)} \in \mathcal{B}^*,$$

moreover

$$u \simeq \Psi_j^{(in)} - \Psi_j^{(out)}, \quad \text{on } \mathcal{M}_j, \quad j = 1, \dots, N,$$

$$u = \Psi_j^{(in)} - \Psi_j^{(out)}, \quad \text{on } \mathcal{M}_j, \quad j = N + 1, \dots, N + N'.$$

We call the mapping

$$\begin{aligned} \mathcal{S}(\lambda) : & (\psi_1^{(in)}, \dots, \psi_N^{(in)}, \{a_\ell^{(N+1)}\}, \dots, \{a_\ell^{(N+N')}\}) \\ & \rightarrow (\psi_1^{(out)}, \dots, \psi_N^{(out)}, \{b_\ell^{(N+1)}\}, \dots, \{b_\ell^{(N+N')}\}) \end{aligned}$$

the **generalized scattering matrix**.

### Theorem (Inverse scattering from cusp)

Let  $\mathcal{M}_{N_1+N'_1}^{(1)}$ ,  $\mathcal{M}_{N_2+N'_2}^{(2)}$  be cusp ends of  $\mathcal{M}^{(1)}$  and  $\mathcal{M}^{(2)}$ . Assume that

- (i)  $\mathcal{M}_{N_1+N'_1}^{(1)}$  and  $\mathcal{M}_{N_2+N'_2}^{(2)}$  are isometric,
- (ii) The components of the **generalized S-matrix** associated with the cusp ends coincide for all energies. Then  $\mathcal{M}^{(1)}$  and  $\mathcal{M}^{(2)}$  are isometric.

## Idea of the proof

Introduce an artificial boundary

$$S = \{(r, x) ; r = 2, x \in M_1\} \subset \mathcal{M}_1.$$

Let

$$\mathcal{M}_{ext} = \mathcal{M}_1 \cap \{r \geq 2\},$$

$$\mathcal{M}_{int} = \mathcal{M} \setminus \mathcal{M}_{ext}.$$

Then  $\widehat{S}_{11}(k)$  determines the N-D map

$$H^1(S) \ni f \rightarrow u|_S,$$

where  $u$  is the solution to the equation

$$(H - \lambda)u = 0, \quad \text{in } \mathcal{M}_{int},$$

$$\frac{\partial u}{\partial r} = f, \quad \text{on } S,$$

$u$  satisfies the radiation condition.

Using this N-D map, one can apply the **boundary control method** due to Belishev-Kurylev to reconstruct the manifold  $\mathcal{M}$ . For the case of cusp, one needs to introduce the generalized S-matrix.

(On the cusp end  $\mathcal{M}_{N+N'}$ , one uses eigenprojections associated with  $-\Delta_h$  to reduce to the 1-dimensional problem, and observe the behavior of growing and decaying solutions.)

## Remarks on the assumption

The decay assumption is roughly as follows. For the sake of simplicity, let us consider only the case of regular ends, and assume that

$$\rho(r)^{-1} \in \mathcal{S}^{-\beta},$$
$$h(r, x, dx) - h_M(x, dx) \in \mathcal{S}^{-\gamma}.$$

Then we assume :




(1) For the limiting absorption principle,

$$\beta > 0, \quad \gamma > 0.$$

(2) For the spectral representation, and the inverse scattering,

$$\text{either } \beta \geq 1, \quad \gamma > 0, \quad \text{or } \beta > 0, \quad \gamma > 1.$$

# For Further Reading I

-  H. Isozaki, Y. Kurylev and M. Lassas,  
Conic singularities, generalized scattering matrix, and  
inverse scattering on asymptotically hyperbolic surfaces,  
to appear in *J. Reine angew. Math.*
-  H. Isozaki, Y. Kurylev and M. Lassas,  
Spectral theory and inverse problems on asymptotically  
hyperbolic orbifolds,  
*arXiv : 1312.0421*
-  H. Isozaki, Y. Kurylev and M. Lassas,  
Recent progress of inverse scattering theory on  
non-compact manifolds,  
to appear in *Contemp. Math.*



# For Further Reading II