# On the integrability of the sine-Gordon equation

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#### Motivation: pseudosphere

 $\Sigma \subset \mathbb{R}^2$  open.  $\exists$  one to one correspondence (Hilbert):

- (i) Immersion  $f: \Sigma \to \mathbb{R}^3, (x,t) \mapsto f(x,t)$  with Gaussian curvature  $\equiv -1$  (pseudosphere) and  $|\partial_x f| = |\partial_t f| = 1$ .
- (ii) solution  $u:\Sigma\to\mathbb{R},(x,t)\mapsto u(x,t)$  to

$$u_{xt} = \sin(u). \tag{1}$$

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" $(i) \Rightarrow (ii)$ " the angle  $\angle(\partial_x f, \partial_y f) =: u(x, y)$  is a solution of (1). " $(ii) \Rightarrow (i)$ "  $\angle(\partial_x, \partial_y) = u(x, y)$  with  $|\partial_x| = |\partial_t| = 1$ characterizes metric on  $\Sigma$ ; integrate Gauss-Mainardi-Codazzi equations to get pseudosphere f.

$$(IVP) \begin{cases} \partial_t \partial_x u(x,t) = \sin(u(x,t)) & t \in \mathbb{R}, \ x \in \mathbb{T} = \mathbb{R}/\mathbb{Z} \\ u(0,x) = u_0(x) & u(x,t) \in \mathbb{T}_{2\pi} = \mathbb{R}/2\pi\mathbb{Z} \end{cases}$$

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 $u \in H^m(\mathbb{T}, \mathbb{T}_{2\pi}), m \ge 1$ :  $\exists \ \ u \in H^m(\mathbb{T}, \mathbb{R}), k \in \mathbb{Z}$  s.t.

$$u(x) = \mathring{u}(x) + 2\pi kx \pmod{2\pi}.$$

**Involution:**  $u \mapsto \breve{u} + \pi, \ \breve{u}(x,t) = u(-x,t)$ 

u solution  $\Leftrightarrow \breve{u} + \pi$  solution

#### The phase space

**observation:** Any  $C^1$  solution  $u : t \mapsto u(t, \cdot) \in H^m(\mathbb{T}, \mathbb{T}_{2\pi})$  satisfies

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Introduce for  $m \ge 1$ :

$$M^{m} := \left\{ u \in H^{m}(\mathbb{T}, \mathbb{T}_{2\pi}) : \int_{\mathbb{T}} e^{iu} \, \mathrm{d}x \in \mathbb{R} \right\}$$
$$M^{m} = M^{m}_{+} \cup M^{m}_{-} \cup M^{m}_{0}$$
$$M^{m}_{0} \subset M^{m} : \int_{\mathbb{T}} e^{iu} \, \mathrm{d}x = 0$$
$$M^{m}_{\pm} \subset M^{m} : \pm \int_{\mathbb{T}} e^{iu} \, \mathrm{d}x \in \mathbb{R}_{>0}$$

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#### The phase space

 $\forall u \in H^m(\mathbb{T}, \mathbb{T}_{2\pi})$  consider the line  $u + \mathbb{T}_{2\pi}$  and note,

$$\int_{\mathbb{T}} e^{\mathbf{i}(u+c)} \, \mathrm{d}x = e^{\mathbf{i}c} \int_{\mathbb{T}} e^{\mathbf{i}u} \, \mathrm{d}x \quad \forall \ c \in \mathbb{T}_{2\pi}.$$

case 1: 
$$\int e^{\mathbf{i}u} dx = 0$$

$$u + c \in M_0^m \quad \forall \ c \in \mathbb{T}_{2\pi}$$

case 2: 
$$\int_{\mathbb{T}} e^{\mathbf{i}u} \, \mathrm{d}x \neq 0$$
$$\exists ! c_{+} \in \mathbb{T}_{2\pi} \text{ s.t. } u + c_{+} \in M^{m}_{+}$$
$$\exists ! c_{-} \in \mathbb{T}_{2\pi} \text{ s.t. } u + c_{-} \in M^{m}_{-}$$
where  $c_{+} = c_{-} + \pi = -\arg\left(\int_{\mathbb{T}} e^{\mathbf{i}u} \, \mathrm{d}x\right) \pmod{2\pi}$ 

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# parametrization of $M^m_+$

Introduce new variable  $v := \frac{1}{2}u_x$ note that  $u \in H^m(\mathbb{T}, \mathbb{T}_{2\pi}) \rightsquigarrow v \in H^{m-1}(\mathbb{T}, \mathbb{R})$ 

#### Introduce for $l \geq 0$

$$W^{l} = \left\{ v \in H^{l}(\mathbb{T}, \mathbb{R}) : \int_{\mathbb{T}} \exp(2\mathrm{i}\partial_{x}^{-1}v) \,\mathrm{d}x \neq 0, \int_{\mathbb{T}} v \,\mathrm{d}x \in \pi\mathbb{Z} \right\}$$

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where  $\partial_x^{-1}$  is the mean zero anti-derivative.

#### Definition

$$\begin{split} \Psi^+ &: W^l \to M_+^{l+1}, \ v \mapsto 2\partial_x^{-1}v + \gamma(v) \\ \Psi^- &: W^l \to M_-^{l+1}, \ v \mapsto \left(2\partial_x^{-1}v\right) \,\check{} + \gamma(v) + \pi \\ \end{split}$$
where  $\check{v}(x) = v(-x) \ and \ \gamma(v) = -\arg\left(\int_{\mathbb{T}} e^{2i\partial_x^{-1}v} \,\mathrm{d}x\right) \,\check{v}$ 

#### Theorem

For any  $l \geq 0$ , the following holds:

(i)  $\Psi^+ [\Psi^-]$  parametrizes  $M^{l+1}_+ [M^{l+1}_-]$  real analytically. Equation (1) in v-coordinates:

$$v_t = \frac{1}{2}\sin(2\partial_x^{-1}v + \gamma(v)) \quad on \ W^l.$$
<sup>(2)</sup>

(ii) Equation (2) is Hamiltonian with respect to the Gardner bracket, i.e. is equivalent to  $v_t = \partial_x \partial_v H$ , with Hamiltonian

$$H: W^{l} \to \mathbb{R}, v \mapsto \frac{1}{4} \int_{\mathbb{T}} \cos(2\partial_{x}^{-1}v + \gamma(v)) \,\mathrm{d}x.$$
(3)

 (iii) The Hamiltonian H is in the Poisson algebra of the focusing mKdV equation and hence (2) is an integrable PDE on W<sup>l</sup>.

#### Corollary

For any  $l \ge 0$ , equation (2) admits Birkhoff coordinates near 0. In particular, near 0, the Hamiltonian H is a function of the actions alone.

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## sketch of the proof

- (i) [KST08]: global real analytic Birkhoff coordinates for defocusing mKdV
- (ii)  $\Rightarrow$  existence of Birkhoff coordinates near 0 for focusing mKdV
- (iii) all Hamiltonians in the Poisson algebra of focusing mKdV, when expressed in Birkhoff coordinates are in normal form.

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$$M_0^m = \left\{ u \in M^m : \int_{\mathbb{T}} \cos(u) \, \mathrm{d}x = 0 \right\}$$

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#### Lemma

 $\forall m \geq 1 \ M_0^m$  is a codimension 2 submanifold of  $H^m(\mathbb{T}, \mathbb{T}_{2\pi})$ 

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#### Definition

For all  $k \in \mathbb{Z}$  introduce

$$M_0^{m,k} = \{ u \in M_0^m) : \exists \ \mathring{u} \in H^m(\mathbb{T}, \mathbb{R}) \ s.t.$$
$$u(x) = \mathring{u}(x) + 2\pi kx \pmod{2\pi} \}$$

Then

$$M_0^m = \bigcup_{k \in \mathbb{Z}} M_0^{m,k} \quad \text{(disjoint union)}$$

#### Theorem

 $\forall \ k \in \mathbb{Z} \ \exists \ U \subset M_0^{1,k} \ open \ s.t. \ \forall \ u_0 \in U, \ \forall \ T > 0 \ the \ initial value \ problem$ 

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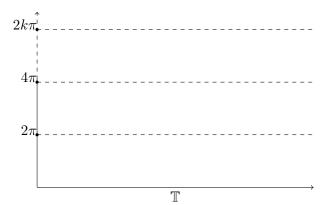
$$\begin{cases} \partial_t \partial_x u = \sin(u) \\ u(0) = u_0 \end{cases}$$

admits no  $C^1$  solution u in  $C^1([0,T], M^1)$ .

#### Remark.

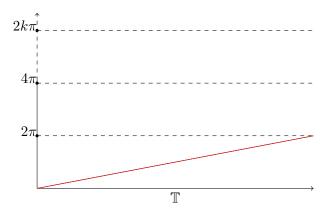
Note that there are elements in U which are  $C^{\infty}$ .

Examples of elements in  $M_0^1$  $(M_0^1 = \{ u \in H^1(\mathbb{T}, \mathbb{T}_{2\pi}), \int_{\mathbb{T}} e^{\mathrm{i}u} \, \mathrm{d}x = 0 \} )$ 



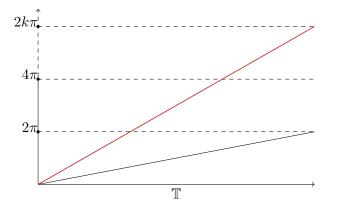
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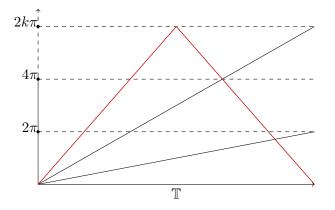
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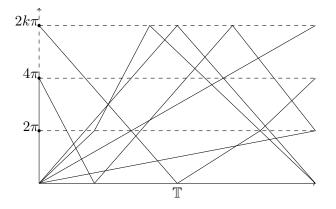
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# $u: t \mapsto u(t, \cdot) \in M^1$ C<sup>1</sup>-smooth solution of (1).

 $u: t \mapsto u(t, \cdot) \in M^1 \ C^1 \text{-smooth solution of } (1).$  $0 = \partial_t \int_{\mathbb{T}} \sin(u) \, \mathrm{d}x = \int_{\mathbb{T}} \cos(u) u_t \, \mathrm{d}x$ 

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 $u_t = \partial_x^{-1} \sin(u) + \partial_t [u]$ , with  $[u] = \int_{\mathbb{T}} u \, dx$  independent of x

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On  $M^m_{\pm}$  (4) is trivially solved when  $u = \Psi^{\pm}(v)$  with v a solution of (2).

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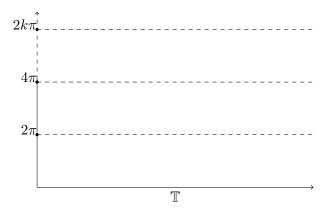
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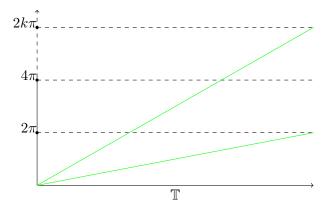
On  $M^m_{\pm}$  (4) is trivially solved when  $u = \Psi^{\pm}(v)$  with v a solution of (2). On  $M^m_0$  (4) reduces to

$$0 = \int_{\mathbb{T}} \cos(u) \partial_x^{-1} \sin(u) \, \mathrm{d}x$$

# When does $0 = \int_{\mathbb{T}} \cos(u) \partial_x^{-1} \sin(u) \, \mathrm{d}x$ hold?

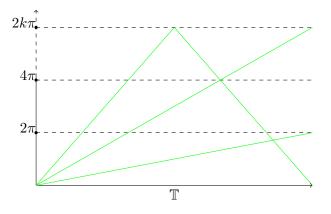


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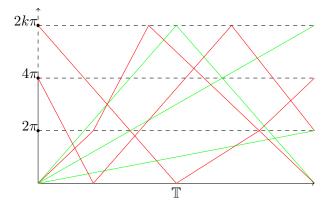
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T Kappeler, B Schaad, and P Topalov, *mKdV and its Birkhoff coordinates*, Physica D: Nonlinear Phenomena **237** (2008), no. 10, 1655–1662.

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