

# Non-uniqueness from molecular equations

St. Petersburg, September 2014  
N.N. Uralceva's 80th birthday

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$$\mu_t^{(N)}(d\vec{x}) = \frac{1}{N} \sum_{i=1}^N \delta_{\vec{x}_i(t)}(d\vec{x}),$$

$$\begin{aligned} M_t^{(N)}(d\vec{x}, d\vec{v}) &= \frac{1}{N} \sum_{i=1}^N \delta_{(x_i, u_i)}(dx, dv) \\ &= (Id \times \vec{u}_N) \# \mu_N(d\vec{x}, d\vec{v}), \end{aligned}$$

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$$\begin{aligned} \frac{d}{dt} \int_{\mathbb{R}^3} \phi(x) u_N(t, x) \mu_t^{(N)}(dx) \\ &= \int_{\mathbb{R}^3} \nabla_x \phi(t, x) \cdot u_N(t, x) u_N(t, x) \mu_t^{(N)}(dx) \\ &\quad - \frac{1}{N\sigma_N} \int \phi(t, x) \Phi' \left( \frac{|x - x'|}{\sigma_N} \right) \frac{x - x'}{|x - x'|} n_t^{(N)}(dx, dx'). \end{aligned}$$

Under assumption on positions and velocities can show:

$$\begin{aligned}\mu_t^{(N)}(d\vec{x}) &\Rightarrow \mu_t(d\vec{x}), \\ M_t^{(N)}(d\vec{x}, d\vec{v}) &\Rightarrow M_t(d\vec{x}, d\vec{v}),\end{aligned}$$

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$$\int \phi(\vec{x}) \vec{u}_N(t, \mathbf{x}) \mu_t^{(N)}(d\vec{x}) \rightarrow \int \phi(\vec{x}) \vec{u}(t, \vec{x}) \mu_t(d\vec{x}), \quad \forall \phi \text{ test.}$$

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$$\int |\vec{u}_N|^2(\vec{x}) \mu^{(N)}(d\vec{x}) \rightarrow \int |\vec{u}|^2(\vec{x}) \mu(d\vec{x}).$$

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weakly.

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Accelerations are uniformly bounded in  $N$ .

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In fact, for the  $M_t$ -measure:

$$\frac{d}{dt} \int \phi(x, v) M_t(dx, dv) = \int \nabla_x \phi(x, v) \cdot v M_t(dx, dv).$$

Can solve:

$$M_t = (S_t)_\# M_0,$$
$$S_t(x, v) = (x + tv, v).$$

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and

$$u(0, x_{ij}) = \begin{cases} (0, 0, 1) & \text{if } i+j \text{ even} \\ (0, 0, -1) & \text{if } i+j \text{ odd,} \end{cases}$$

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Then

$$M_0^{(N)}(dx, dv)$$

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$$M_0(dx, dv) = \Delta_0(dx) \otimes \left( \frac{1}{2} \delta_{(0,0,1)}(dv) + \frac{1}{2} \delta_{(0,0,-1)}(dv) \right)$$

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Therefore: initial macroscopical velocity

$$u(0, x) = 0.$$

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 \end{aligned}$$

It gives the macroscopic velocity for  $t > 0$ :

$$u(t, x) = \begin{cases} (0, 0, 1) & (x_1, x_2) \in D, \quad x_3 = t \\ (0, 0, -1) & (x_1, x_2) \in D, \quad x_3 = -t \\ 0 & \text{otherwise.} \end{cases}$$

Not even local uniqueness for

$$\partial_t(u\mu_t) + \operatorname{div}(u \otimes u \mu_t) = 0$$

$$\partial_t\mu_t + \operatorname{div}(u \mu_t) = 0.$$

( $\mu_t$  not absolutely continuous on  $\mathbb{R}^3$ .)

Compatible with macroscopic definition of initial condition for weak solutions:

$$\int \phi(x)u(t, x)\mu_t(dx) - \int \phi(x)u(0, x)\mu_0(dx) = \int_0^t \int \nabla\phi(x) \cdot u(s, x)u(s, x)\mu_s(dx) ds$$

given that

$$\lim_{t \rightarrow 0} \int \phi(x)u(t, x)\mu_t(dx) = \int \phi(x)u(0, x)\mu_0(dx).$$



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Define

$$u_N(x_{ijk}) = \begin{cases} (0, 0, c) & i+j \text{ even} & k \text{ even} \\ (0, 0, -c) & i+j \text{ odd} & k \text{ even} \\ (0, 0, -c) & i+j \text{ even} & k \text{ odd} \\ (0, 0, c) & i+j \text{ odd} & k \text{ odd.} \end{cases}$$

and denote the corresponding measure in  $\mathbb{R}^6$  with  $M^N(dx, dv)$ .

As  $N \rightarrow \infty$ ,

$$M^N(dx, dv)$$

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$$M_0(dx, dv) = \chi_K(x) dx \otimes \frac{1}{2} (\delta_{(0,0,c)}(dv) + \delta_{(0,0,-c)}(dv)).$$

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Macroscopic initial conditions:

$$\begin{cases} \rho(0, x) = \chi_K(x) \\ u(0, x) = 0. \end{cases}$$

$$M_t(dx, dv) =$$

$$\left\{ \begin{array}{ll} \frac{1}{2} dx \otimes \delta_{(0,0,-c)}(dv) & 0 \leq t \leq \frac{l}{c} \quad -l - ct \leq x_3 \leq -l + ct \\ \frac{1}{2} dx \otimes \delta_{(0,0,-c)}(dv) \\ \quad + \frac{1}{2} dx \otimes \delta_{(0,0,c)}(dv) & 0 \leq t \leq \frac{l}{c} \quad -l + ct \leq x_3 \leq l - ct \\ \frac{1}{2} dx \otimes \delta_{(0,0,c)}(dv) & 0 \leq t \leq \frac{l}{c} \quad l - ct \leq x_3 \leq l + ct \\ \frac{1}{2} dx \otimes \delta_{(0,0,-c)}(dv) & t \geq \frac{l}{c} \quad -l - ct \leq x_3 \leq l - ct \\ \frac{1}{2} dx \otimes \delta_{(0,0,c)}(dv) & t \geq \frac{l}{c} \quad -l + ct \leq x_3 \leq l + ct \end{array} \right.$$

Macroscopic density and velocity are

$$\rho(t, x) = \left\{ \begin{array}{ll} \frac{1}{2} & 0 \leq t \leq \frac{l}{c} \quad -l - ct \leq x_3 \leq -l + ct \\ 1 & 0 \leq t \leq \frac{l}{c} \quad -l + ct \leq x_3 \leq l - ct \\ \frac{1}{2} & 0 \leq t \leq \frac{l}{c} \quad l - ct \leq x_3 \leq l + ct \\ \frac{1}{2} & t \geq \frac{l}{c} \quad -l - ct \leq x_3 \leq l - ct \\ \frac{1}{2} & t \geq \frac{l}{c} \quad -l + ct \leq x_3 \leq l + ct, \end{array} \right.$$



$$u(t, x) = \begin{cases} (0, 0, -c) & 0 \leq t \leq \frac{l}{c} & -l - ct \leq x_3 \leq -l + ct \\ 0 & 0 \leq t \leq \frac{l}{c} & -l + ct \leq x_3 \leq l - ct \\ (0, 0, c) & 0 \leq t \leq \frac{l}{c} & l - ct \leq x_3 \leq l + ct \\ (0, 0, -c) & t \geq \frac{l}{c} & -l - ct \leq x_3 \leq l - ct \\ (0, 0, c) & t \geq \frac{l}{c} & -l + ct \leq x_3 \leq l + ct. \end{cases}$$

Fluctuations last within time interval  $0 \leq t \leq \frac{l}{c}$ .

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Equations satisfied weakly:

$$\partial_t \rho(t, \mathbf{x}) + \operatorname{div}(\rho(t, \mathbf{x}) \mathbf{u}(t, \mathbf{x})) = 0$$

$$\begin{aligned} \partial_t(\rho(t, \mathbf{x}) \mathbf{u}(t, \mathbf{x})) + \operatorname{div}(\rho(t, \mathbf{x}) \mathbf{u}(t, \mathbf{x}) \otimes \mathbf{u}(t, \mathbf{x})) = \\ - \underbrace{c^2(0, 0, \partial_3 \chi_{B_t}(\mathbf{x}))}_{\text{from averages of fluctuations}} \end{aligned}$$

## Example Approximate by molecules initial distribution

$$M_0(dx, dv) = C\rho(x) \left\{ \begin{array}{ll} e^{-v^{1/4}} & v > 0 \\ 0 & v \leq 0 \end{array} \right\} dx dv$$

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Approximation in such way that same evolution equation for  $M_t$  holds.

Initial densities, velocities, and all moments: equal at  $t = 0$ .

Flow under equation for measure  $M_t$ .

Already velocities at some  $t > 0$  differ.

## Example Approximate by molecules initial distribution

$$M_0(dx, dv) = C\rho(x) \left\{ \begin{array}{ll} e^{-v^{1/4}} & v > 0 \\ 0 & v \leq 0 \end{array} \right\} dx dv$$

and approximate by molecules the same after changing the velocity distribution

$$M_0(dx, dv) = C\rho(x) \left\{ \begin{array}{ll} e^{-v^{1/4}} + \varepsilon e^{-v^{1/4}} \sin v^{1/4} & v > 0 \\ 0 & v \leq 0 \end{array} \right\} dx dv$$

Approximation in such way that same evolution equation for  $M_t$  holds.

Initial densities, velocities, and all moments: equal at  $t = 0$ .

Flow under equation for measure  $M_t$ .

Already velocities at some  $t > 0$  differ.

Infinitely many macroscopic equations.